

(6) 超新星爆発での 重元素 (Rプロセス元素) 合成と 銀河年齢

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超新星爆発と重元素の起源

重元素＝鉄よりも重い質量を持つトリウムやウランに至る元素

宇宙論との関わり 「宇宙年齢」

宇宙論パラメータによる推定。137 +/- 2 億年。 --- Model dependent!

∴ We don't know the true nature of DARK MATTER nor DARK ENERGY.

核宇宙年代学 (Nucleo-cosmochronology)

^{232}Th (半減期140.5億年)、 ^{238}U (半減期45.7億年)を初期世代天体に検出し、初期天体の年齢を推定。これにより、銀河・宇宙年齢の下限値を推定する。

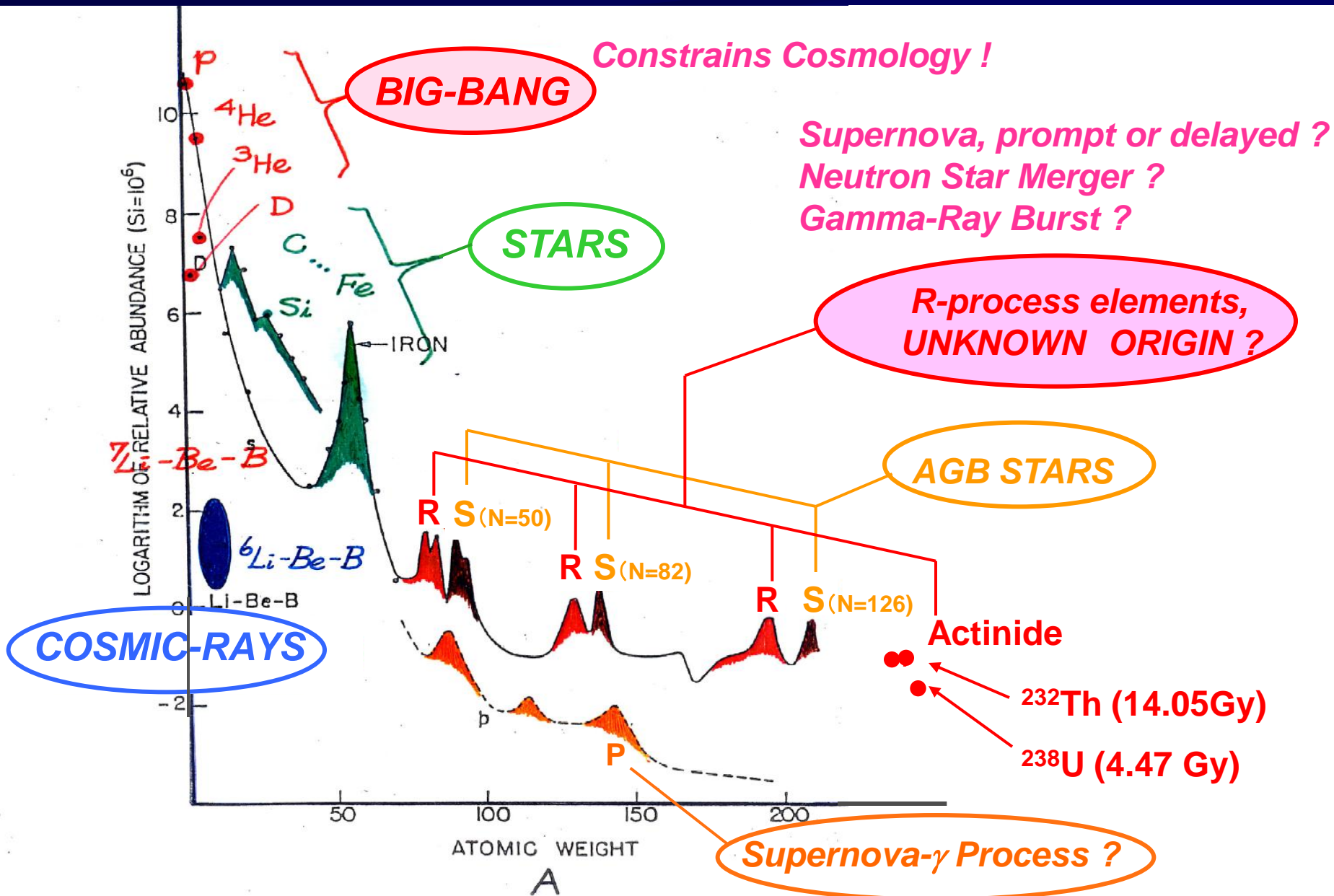
宇宙論パラメータにまったく依存しない。

初期元素量を推定する必要がある。

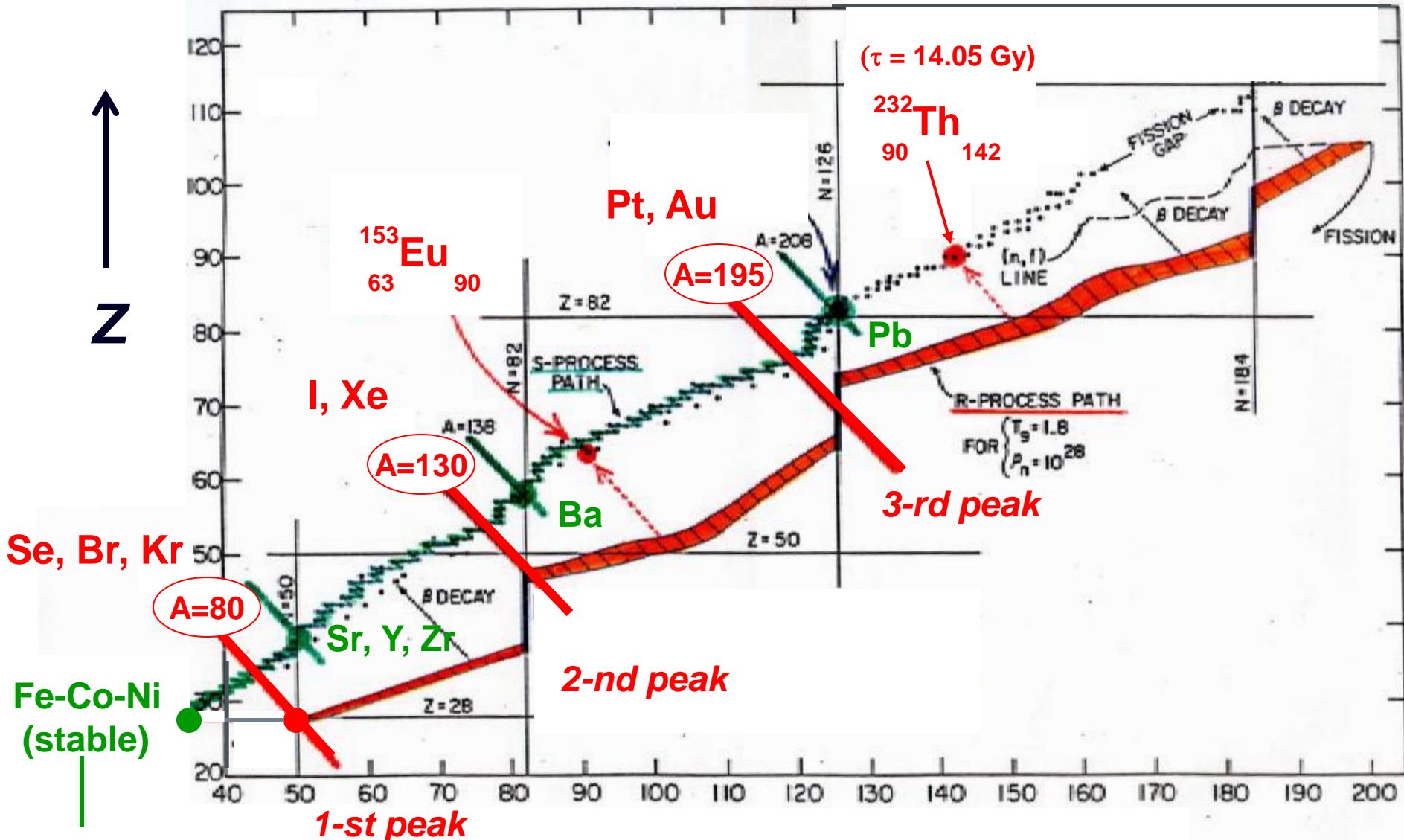
超新星爆発、重元素合成モデル、ニュートリノ相互作用の精密化。

宇宙の化学進化、重元素の起源論との関わり。

Solar System Abundance



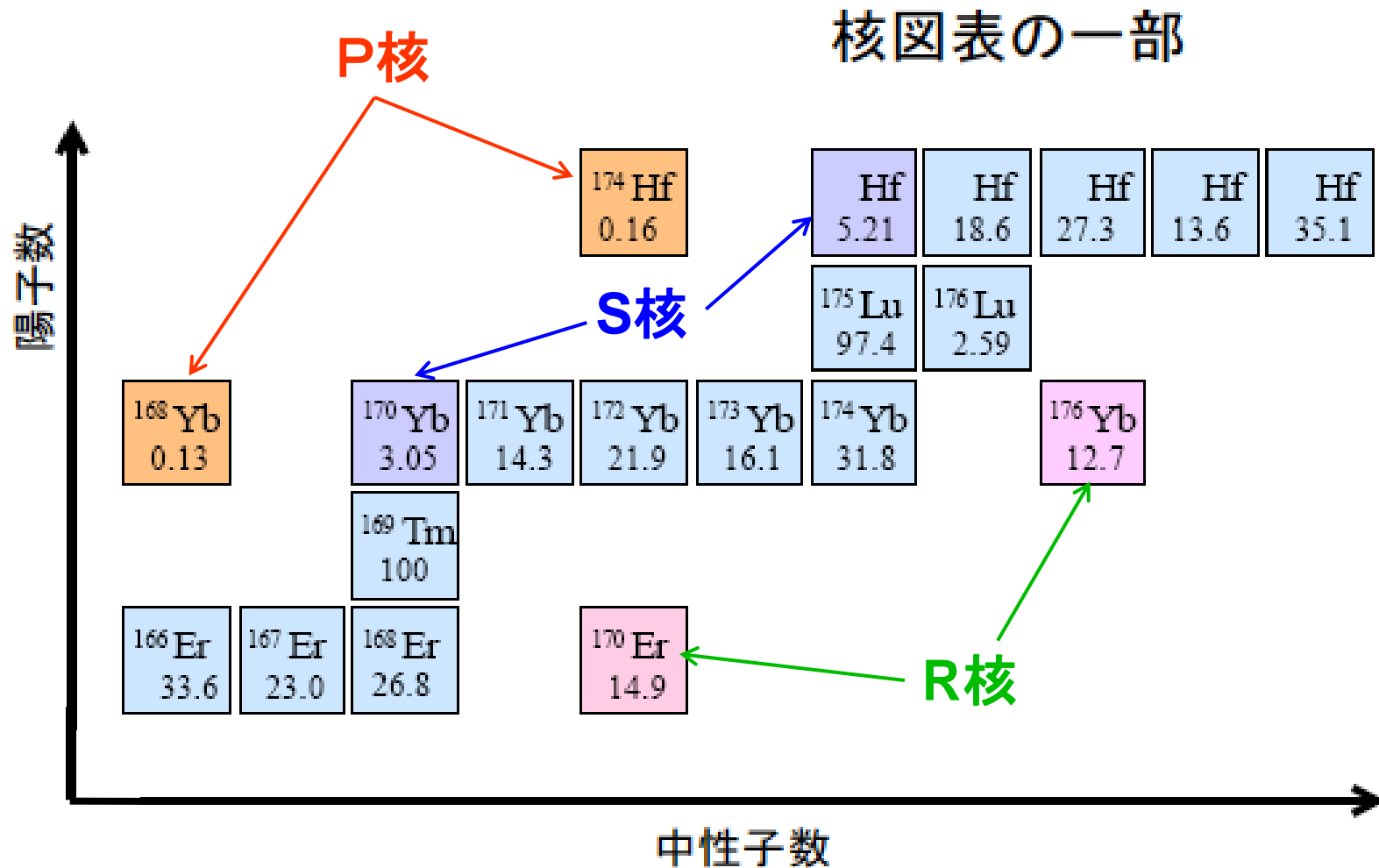
Very Rapid Neutron-Capture Process



These are not "seeds" because r-process is a primary process starting from protons & neutrons !

N →

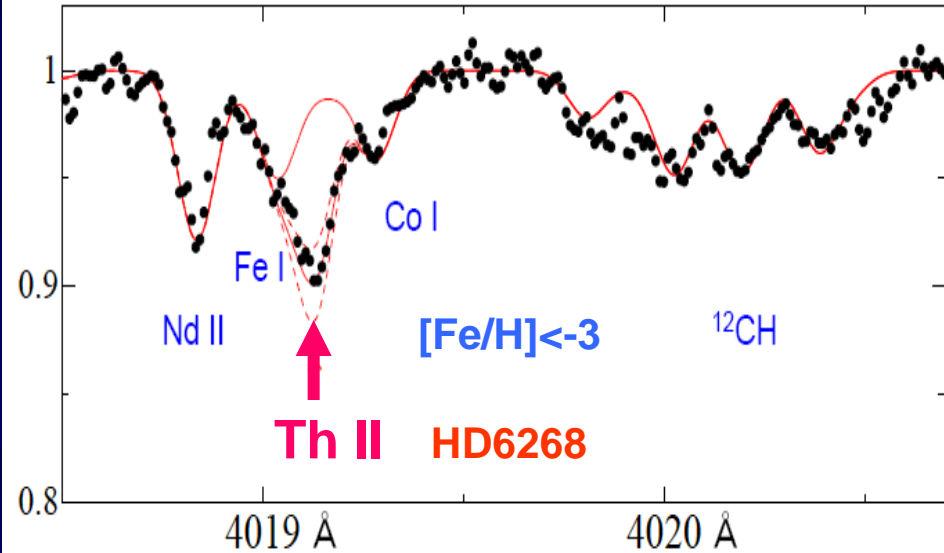
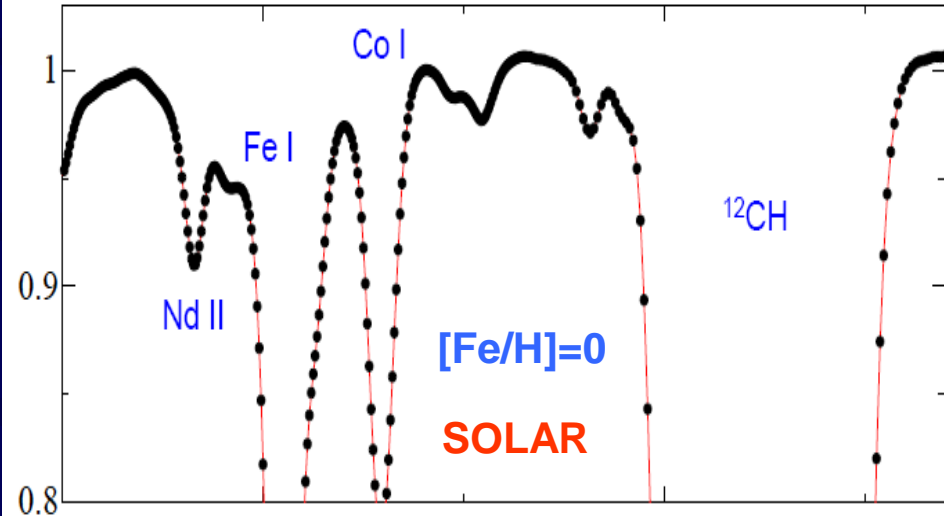
核図表と重元素の起源



- 鉄より重い重元素の約99%は、中性子捕獲反応で生成された。
- その一方で、p核と呼ばれる同位体の起源は謎であった。

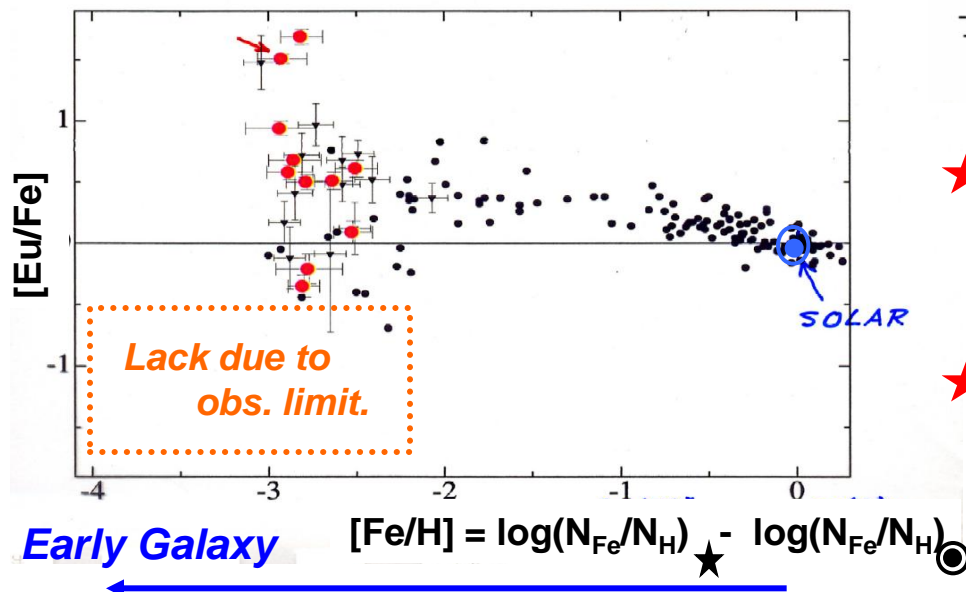
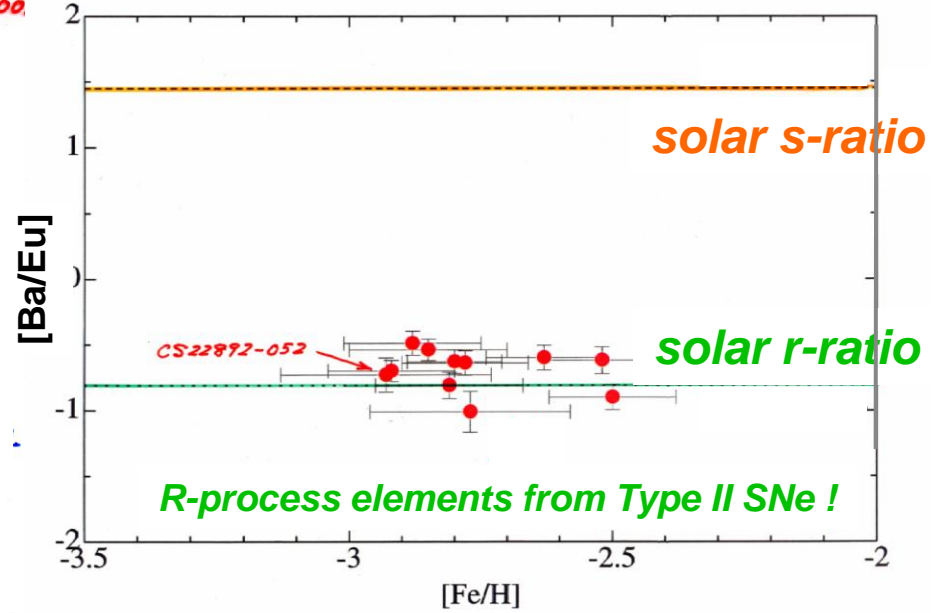
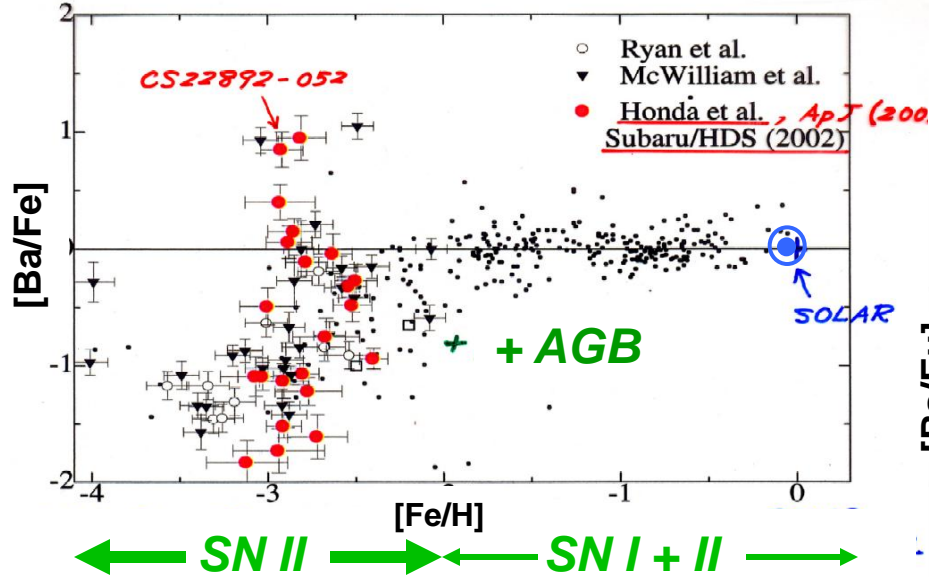
Subaru Telescope

OBSEVES Extremely Metal-Deficient Stars



SUBARU Telescope HDS

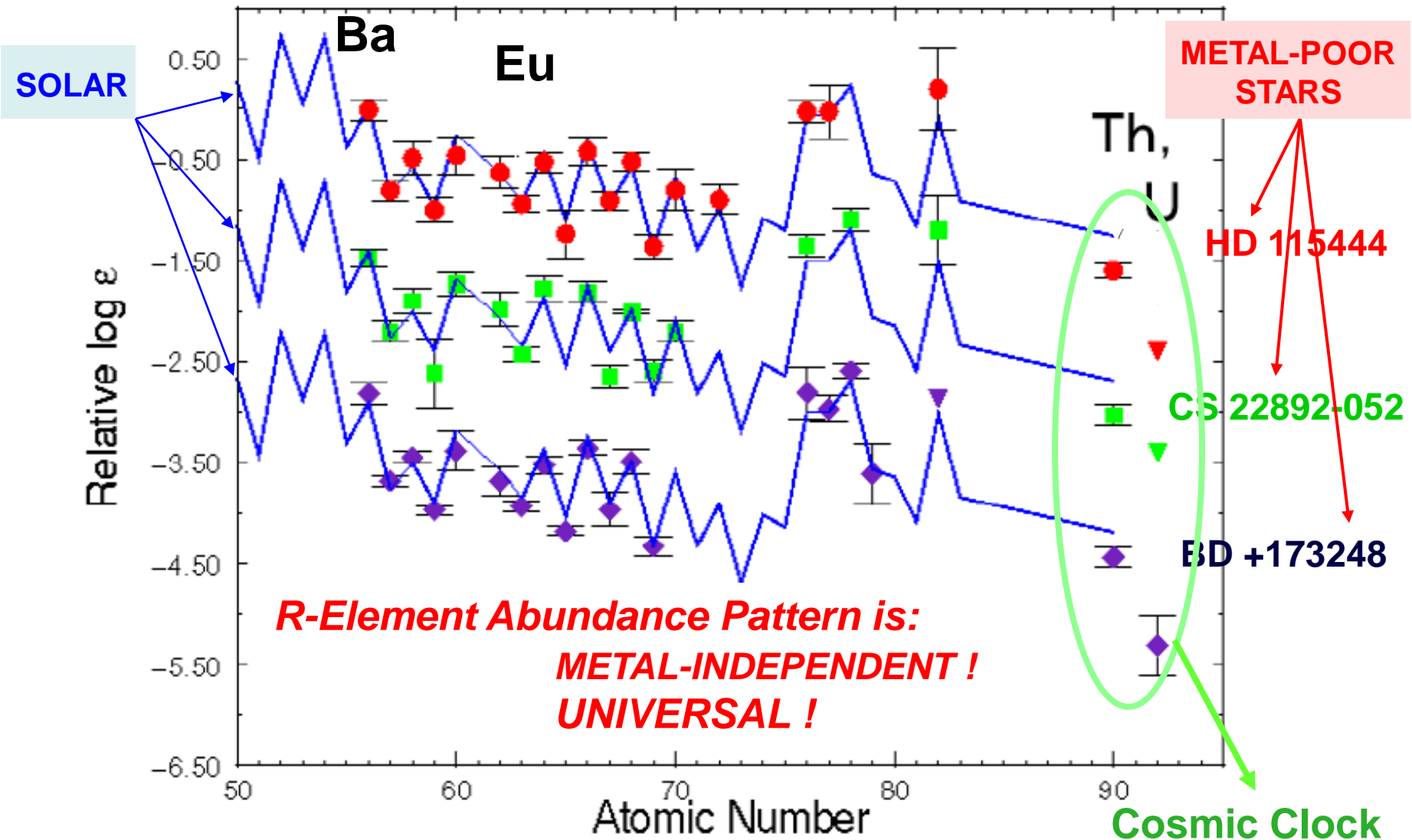
Honda, Aoki, Kajino et al.
 (SUBARU/HDS Collaboration),
 2004, ApJS 152, 113; 2004, ApJ 607, 474



- ★ Large abundance scatter at $[Fe/H] < -2$ is an evidence for INDIVIDUAL supernova episode.
- ★ Only Core-Collapse TYPE II SUPERNOVAE are the likely astrophysical sites of the R-Process !

UNIVERSAL SCALING OF R-PROCESS ABUNDANCES

C. Sneden et al. (1996 – 2005)



Collapse of the Core

Prompt core bounce

$$E(\text{iron core}) \sim GM^2/r \sim 10^{51} \text{ erg}$$

$$E(\text{neutron star}) \sim GM^2/r \sim 10^{53} \text{ erg}$$

$$E(\text{neutron star}) - E(\text{iron core}) \sim 10^{53} \text{ erg}$$

99% is emitted as neutrinos!

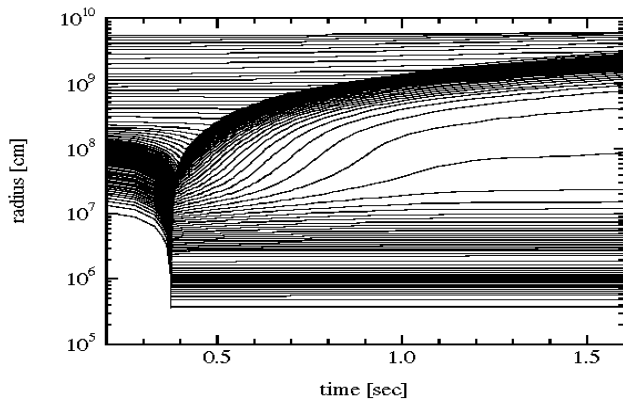
$$E(\text{shock}) \sim 10^{51} \text{ erg}$$

1% is kinetic energy!

Usually the shock is absorbed by dissociating the iron core.

Neutrino-heated explosion

DELAYED SUPERNOVA ⁹



Steps to a Core Collapse Supernova

- Stars with $M \sim 10 - 40 M_{\odot}$ build up an Fe/Ni core.
Maximum core size $M_{\text{ch}} = 5 Y_e^2 M_{\odot} \sim 1.3 M_{\odot}$ (Electron Capture).
- Collapse Separates,
inner homologous ($v \propto r$) core = $1.1 M_{\odot}$.
outer slowly collapsing core = $0.2 M_{\odot}$.

The central density increases and reaches nuclear matter density,
 $\rho_{\text{nucl}} \sim 2 \times 10^{14} \text{ g cm}^{-3}$ (Nuclear EOS).

- An outward moving shock develops due to nuclear saturation.
- The shock dissociates the outer iron core into free nucleons.
- Neutrinos scatter off the heated material behind the shock and deposit energy into p, n, and e^+e^- .
- A high entropy heated region forms and begins to lift the outer layers of the star (neutrino-driven wind).

We detected ν 's, then NEUTRON STAR once formed !



SN1987A

Can core-collapse supernova produce R-PROCESS elements like ^{232}Th ($\tau_{1/2}=14.05\text{Gy}$) which is an celestial cosmic clock ?

QuickTime™ and a
YUV420 codec decompressor
are needed to see this picture.

→ 10 km

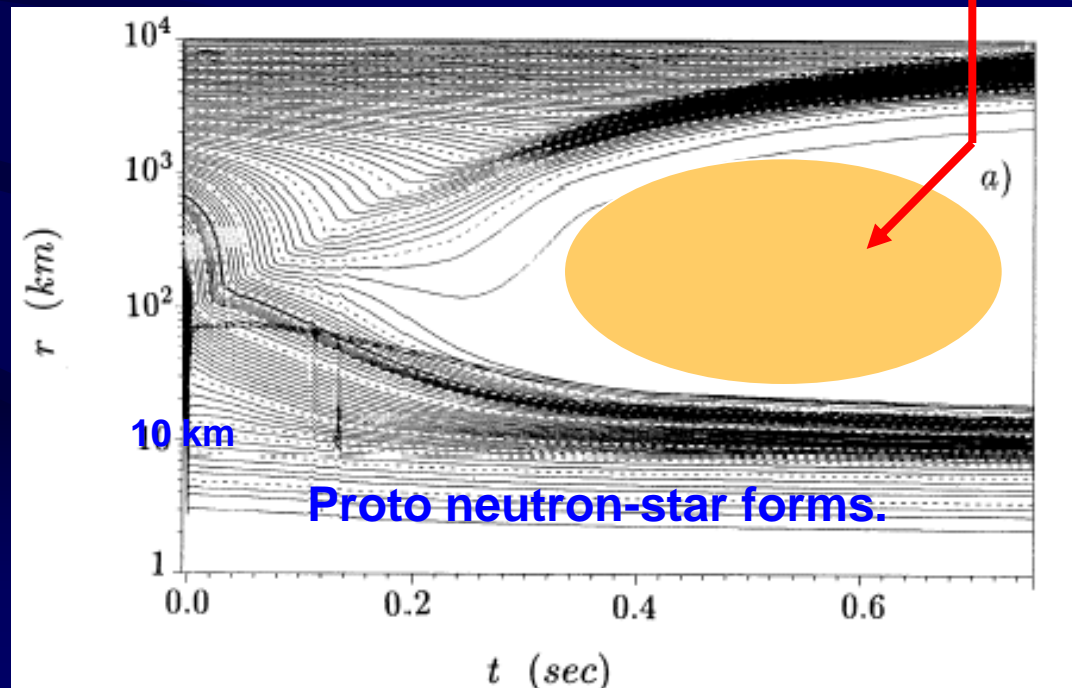
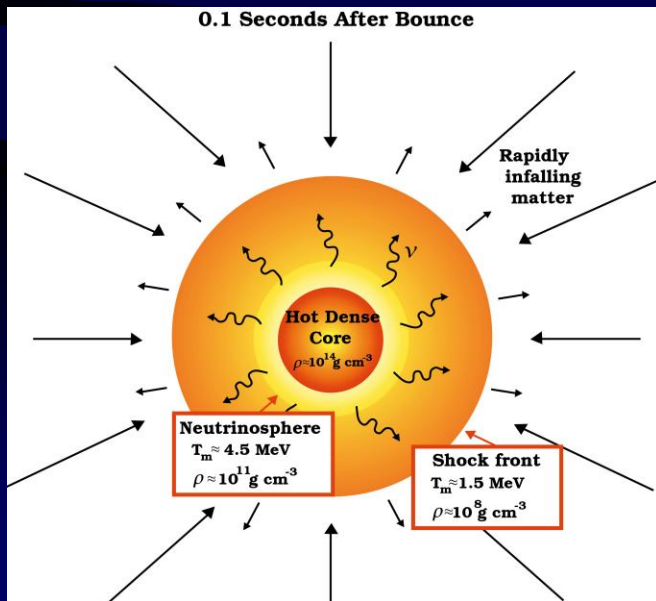
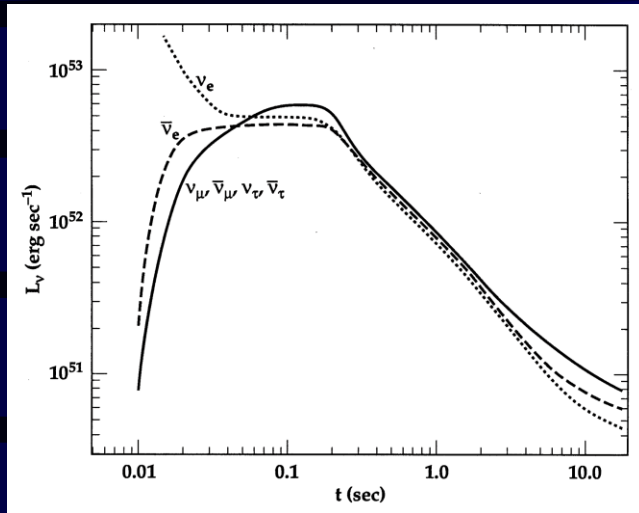
300 km

Neutrino Heated Bubble forms

Neutrino
Luminosity
 $\sim 10^{53}$ erg/s

Neutrino Heating
produces
a high entropy bubble.

Woosley et al. 1994, ApJ 433, 229



General Relativistic Models of ν -Driven Winds

Otsuki, Tagoshi, Kajino and Wanajo 2000, ApJ 533, 424

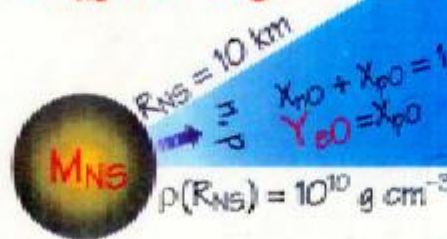
spherically symmetric, steady state winds in Schwarzschild geometry.

$$\left\{ \begin{array}{l} \dot{M} = 4\pi r^2 \rho_b u = \text{const} \quad \text{: mass ejection rate} \\ u \frac{du}{dr} = \frac{1}{\rho_{\text{tot}} + P} \frac{dP}{dr} \left(1 + u^2 - \frac{2M}{r} \right) - \frac{M}{r^2} \quad \text{: equation of motion} \\ \dot{q} = u \left(\frac{d\varepsilon}{dr} - \frac{P}{\rho_b^2} \frac{d\rho_b}{dr} \right) \quad \text{: heating rate} \end{array} \right.$$

$$M = \frac{\text{heating rate}}{\text{grav. potential}}$$

S/k \Rightarrow increase entropy $\sim 200 N_A k - 140 N_A k$
(factor of ~ 2 for $M_{\text{NS}} \sim 2.0 M_{\odot}$)

τ_{exp} \Rightarrow reduce dynamical timescale
(factor of ~ 2 for $M_{\text{NS}} \sim 2.0 M_{\odot}$)



Nucleosynthesis + Diffusion Equation for $Z < 100$ (~3000 species)

Given T & ρ :

$$\frac{dn_A}{dt} = - \sum_{jkl} \langle \sigma v \rangle_{A_j \rightarrow k l} n_A n_j$$

Nuclear
Reactions

$$+ \sum_{kl} \langle \sigma v \rangle_{kl \rightarrow A_j} n_k n_l$$

$$+ (\text{THREE BODY \& HIGHER TERMS})$$

$$- \frac{\ln 2}{T_{1/2}(\beta)} n_A$$

β -decays

$$+ (\text{OTHERS})$$

$$+ \vec{\nabla} \cdot \vec{D}_A \vec{\nabla} n_A$$

Diffusion

Thermonuclear Reaction Rate

$$\tau_i^{-1} \equiv \rho_B N_A \langle \sigma v \rangle_{ij \rightarrow k l}$$

Boltzmann average

$$= \rho_B N_A \sqrt{\frac{8}{\mu \pi}} \frac{1}{(kT)^{3/2}} \int_0^{\infty} E \sigma_{ij \rightarrow k l}(E) \exp(-E/kT) dE$$

Cross Section

SUPERNOVA R-PROCESS

Otsuki, Tagoshi, Kajino & Wanajo

2000, ApJ 533, 424

Wanajo, Kajino, Mathews & Otsuki

2001, ApJ 554, 578

$t = 0$

Neutrino-driven wind forms
right after SN core collapse.

$n + p + \alpha$

$t = 18 \text{ ms}$

Seeds form.

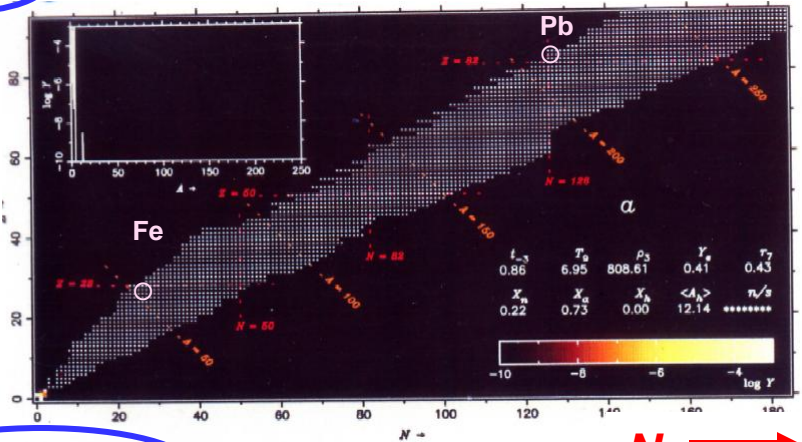
Exotic neutron-rich (^{78}Ni)

$t = 568 \text{ ms} - 1 \text{ s}$

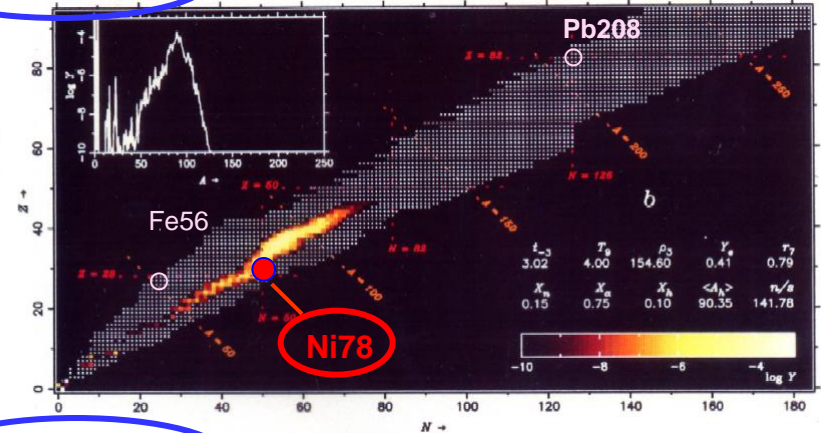
Heavy r-elements synthesize.

$t = 0$

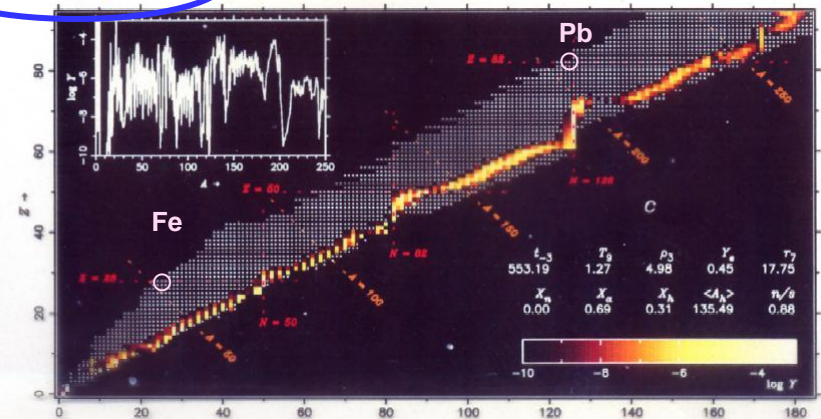
Z



$t = 18 \text{ ms}$



$t = 568 \text{ ms}$



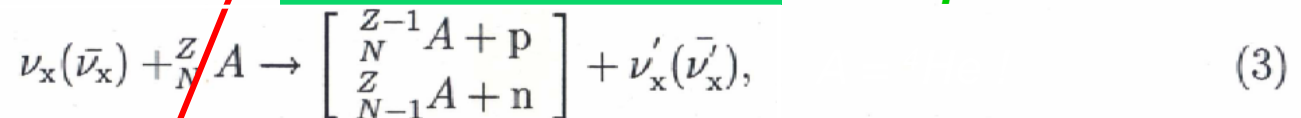
Neutrino Effects on Black Hole vs. Neutron Star Formation

Sasaqui, Kajino & Balantekin 2005, ApJ, in press. (astro-ph/0506100)

The important neutrino reactions during the nucleosynthesis are



Pauli Blocking →



where $x = \mu, \tau$ are the neutrino flavors, and ${}^Z_N A$ is the nucleus with proton number Z and neutron number N . In particular the charged-current reactions that determine the initial neutron-to-proton ratio are

Early ν -cutoff

\therefore makes them impotent!

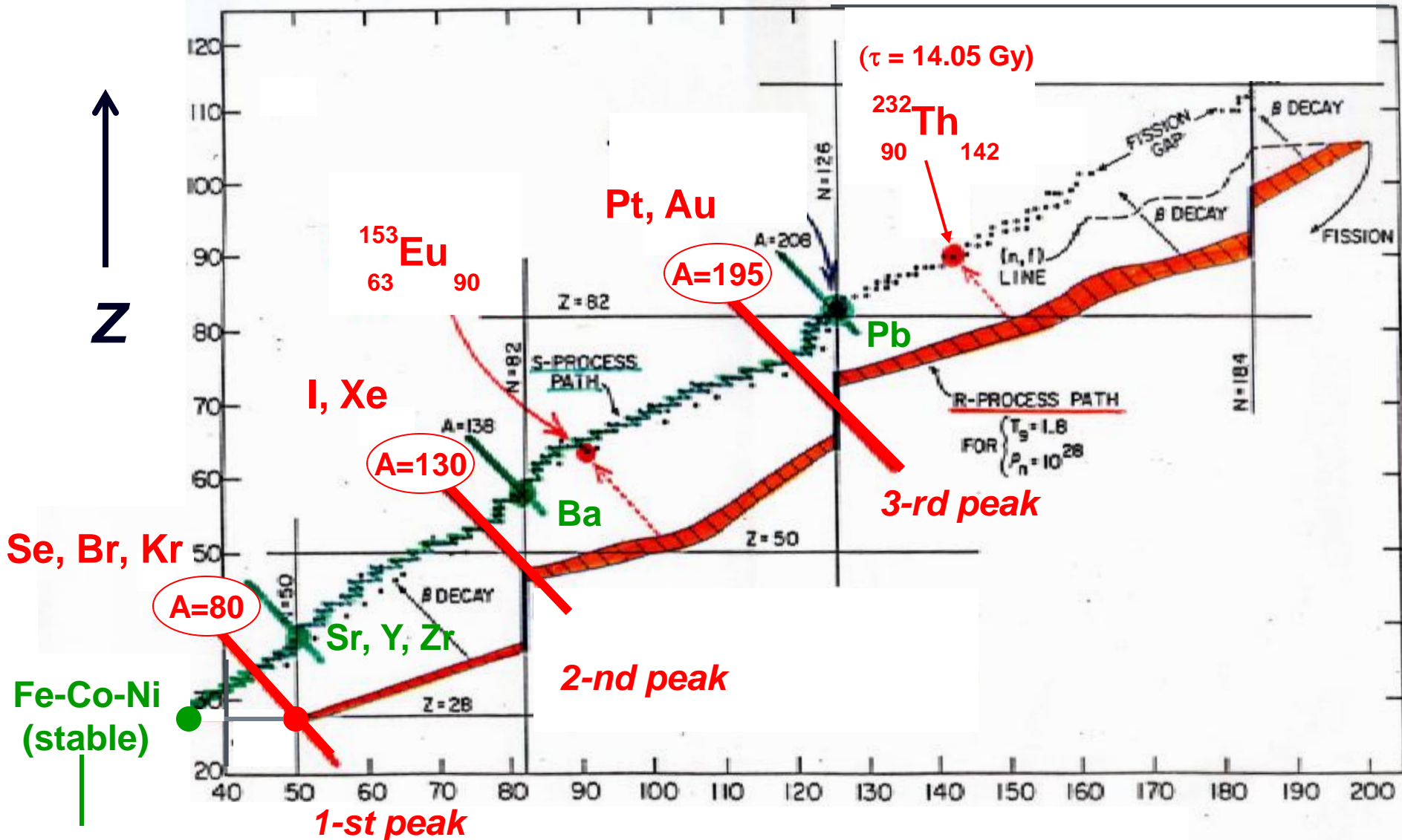
\therefore keeps neutron-rich!



The neutron to proton ratio in the weak equilibrium satisfies (Qian & Woosley 1996),

$$Y_e = \frac{p}{n+p} \approx \left(1 + \frac{L_{\bar{\nu}_e}}{L_{\nu_e}} \times \frac{\epsilon_{\bar{\nu}_e} - 2\Delta + 1.2\Delta^2/\epsilon_{\bar{\nu}_e}}{\epsilon_{\nu_e} + 2\Delta + 1.2\Delta^2/\epsilon_{\nu_e}} \right)^{-1} < 0.5$$

Very Rapid Neutron-Capture Process



These are not "seeds" because r-process is a primary process starting from protons & neutrons !

$N \longrightarrow$

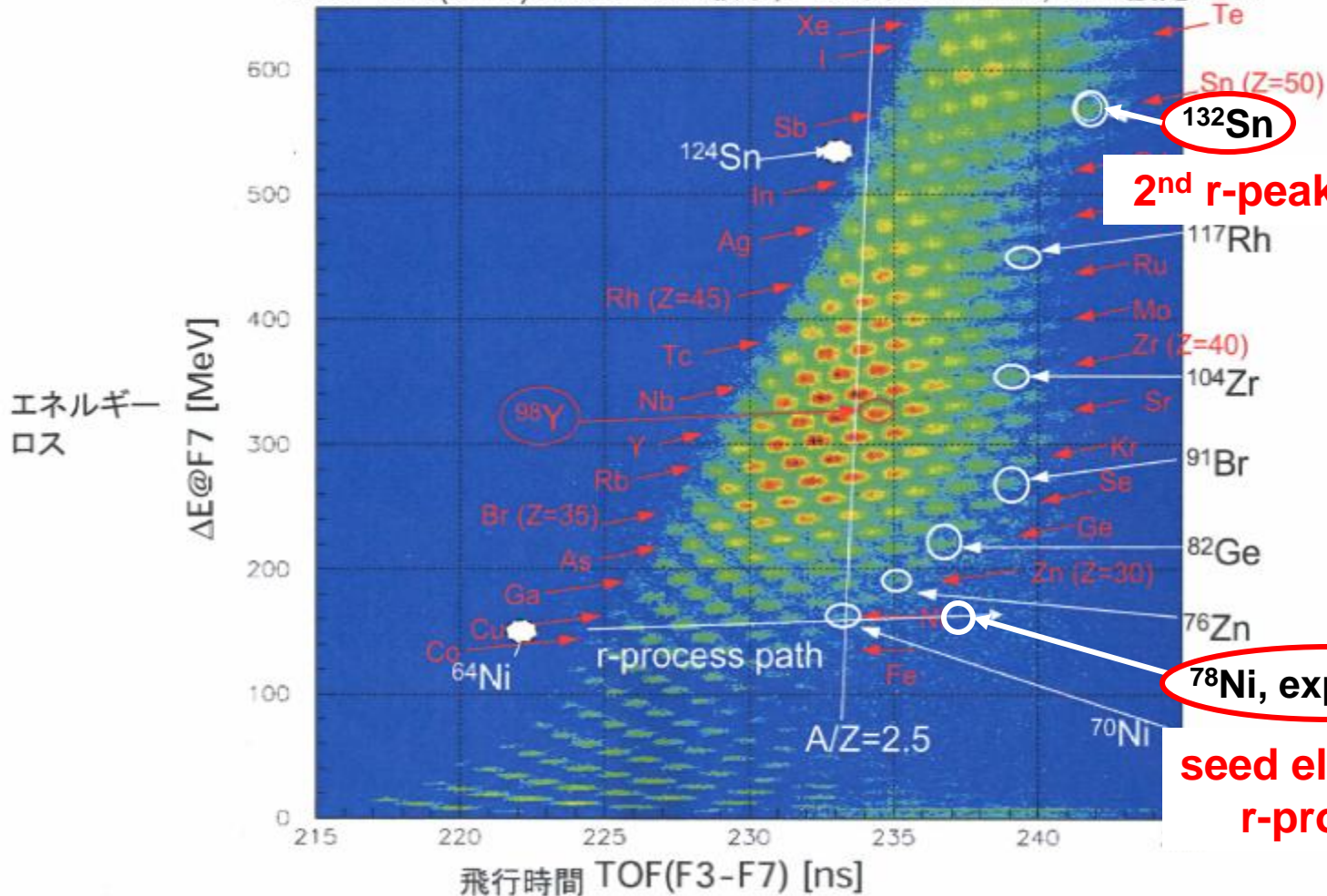
RIKEN-RIBF New Ring Cyclotron (2007)

理研 久保敏幸氏より

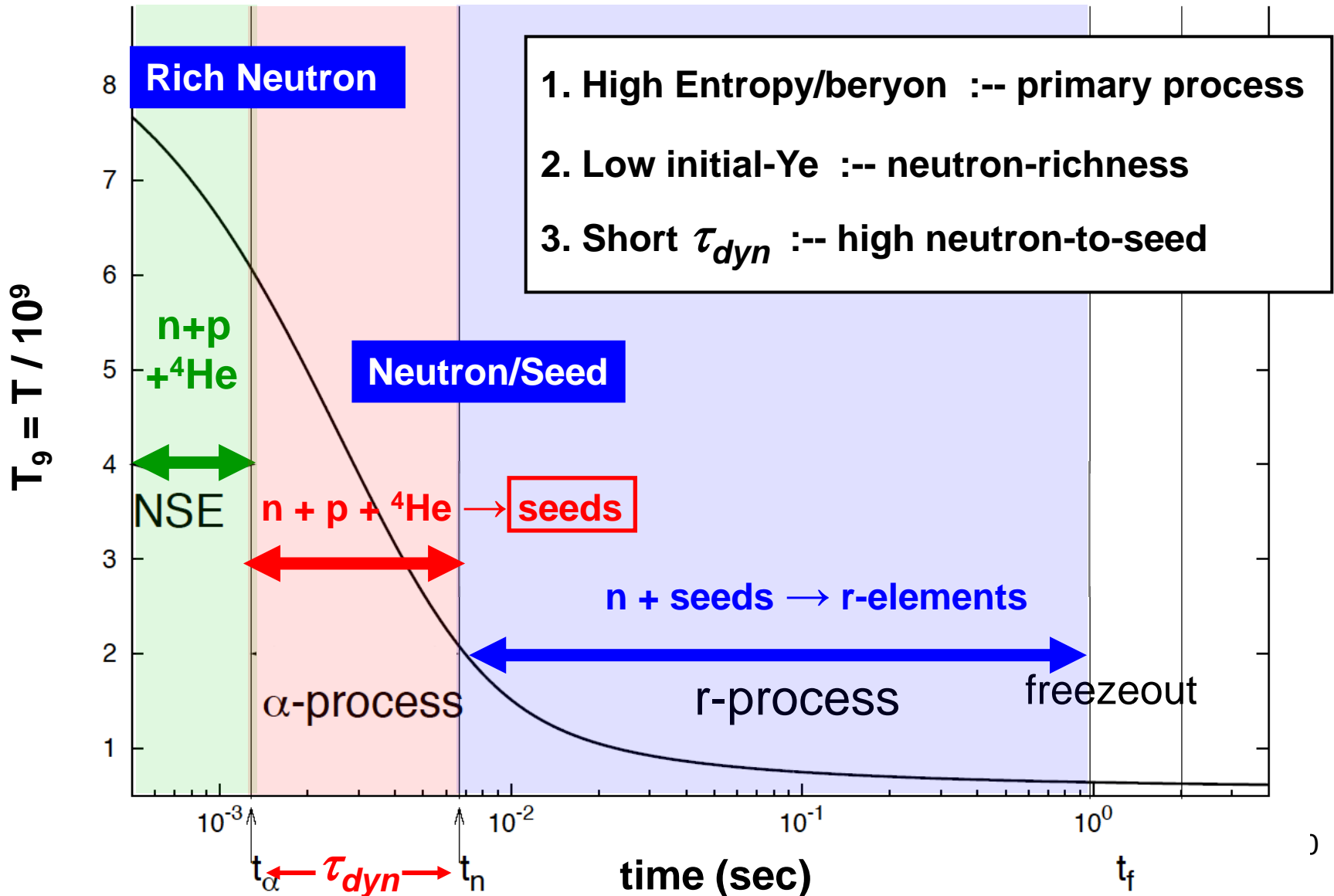
2007年3月26日-27日(測定)

粒子の同定(粒子識別図、PID図): F1デグレーダー無し

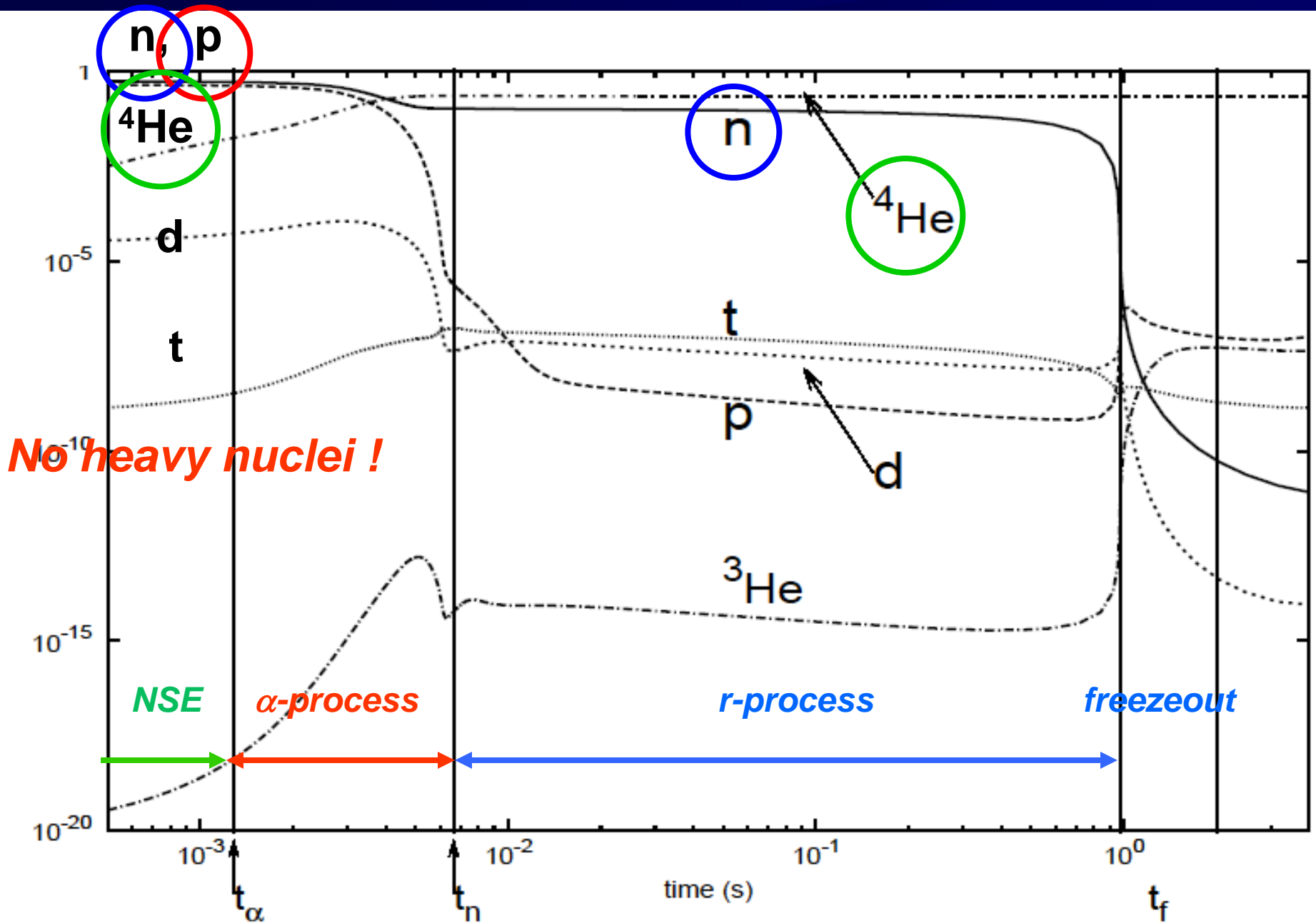
$^{238}\text{U} + \text{Be}(5\text{mm})$ at 345 MeV/核子, F1スリット: +2mm, Brho設定: ^{76}Ni



Nucleosynthesis proceeds: *NSE* → *α*-process → *r*-process



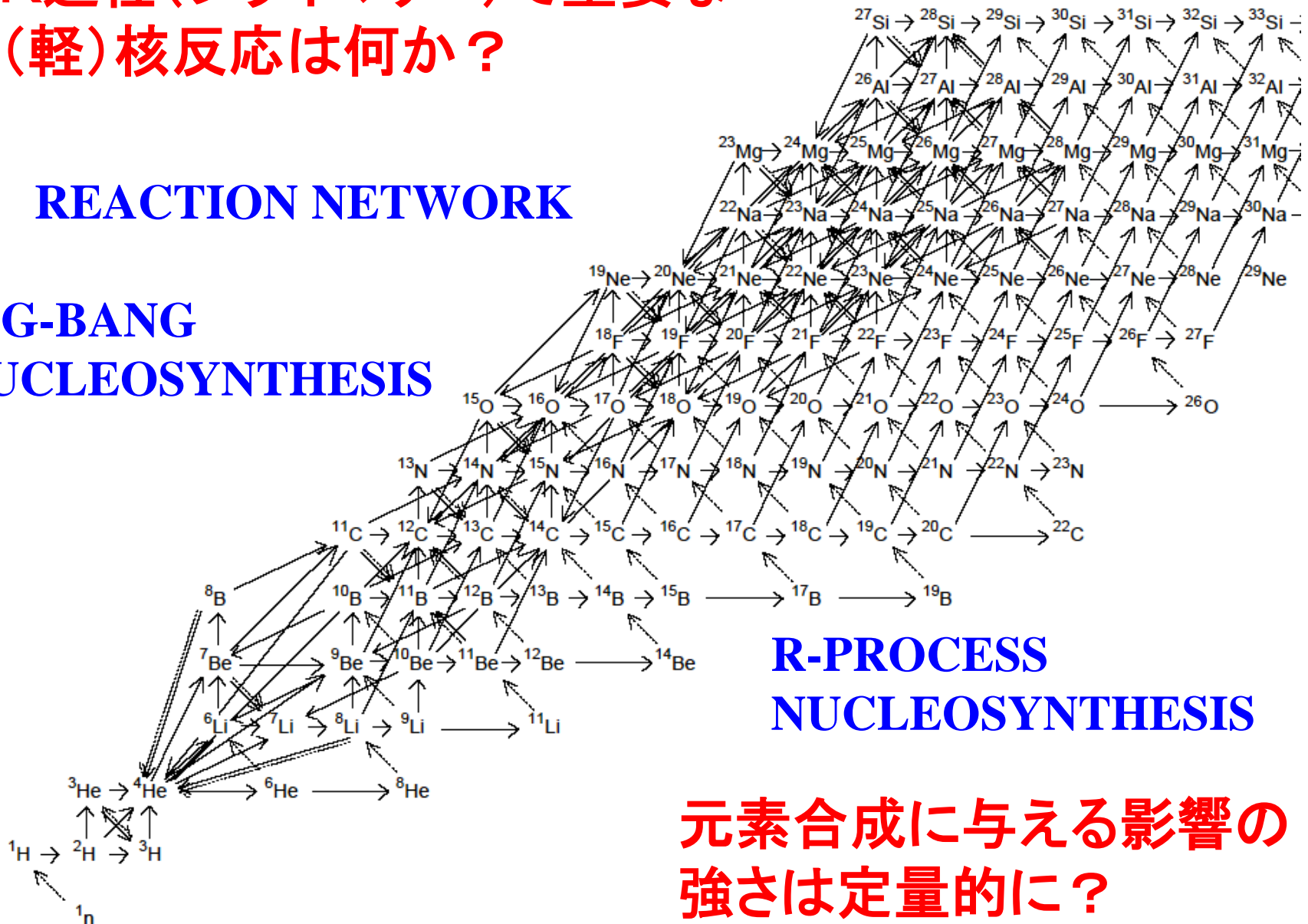
Nucleosynthesis in SN ν - Driven Wind



R過程(プライマリー)で重要な
(軽)核反応は何か？

REACTION NETWORK

BIG-BANG NUCLEOSYNTHESIS



元素合成に与える影響の
強さは定量的に？

Reaction Sensitivity

Sasaqui, Kajino, Mathews, Otsuki & Nakamura
Astrophys. J. (2005), submitted.

$$Y_{2\text{nd}} = Y_{2\text{nd}}(0) \prod_i \left(\frac{S_i}{S_i(0)} \right)^{\sigma_i}$$

$$Y_{3\text{rd}} = Y_{3\text{rd}}(0) \prod_i \left(\frac{S_i}{S_i(0)} \right)^{\sigma_i}$$

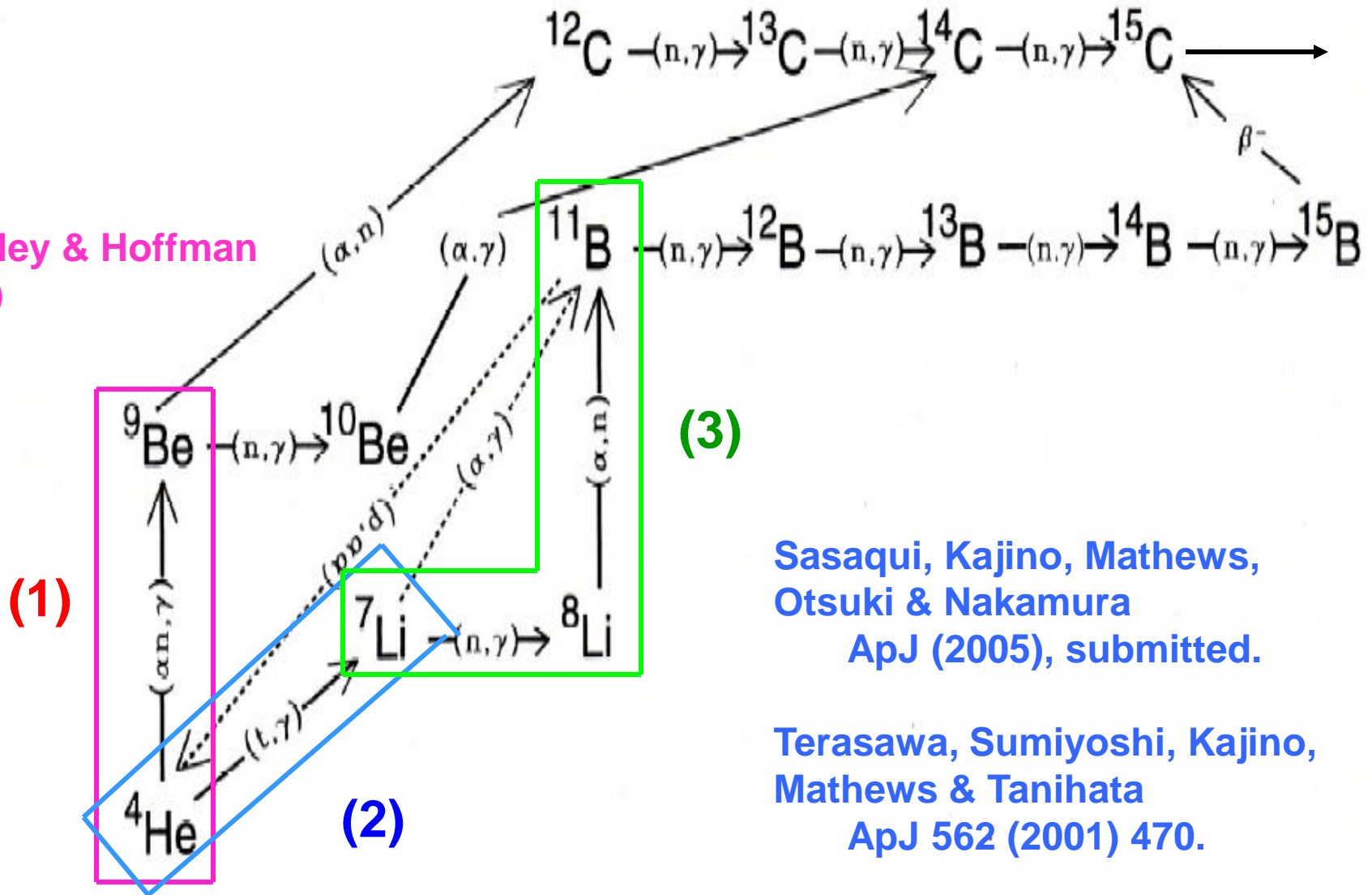
$$\sigma_i = \frac{\partial \left(\log \frac{Y_j}{Y_j(0)} \right)}{\partial \left(\log \frac{S_i}{S_i(0)} \right)}$$

Solar Neutrino Flux (J. Bahcall, Rev. Mod. Phys. 1982)

$$R = 1.35 \text{SNU} \times \left(\frac{S_{11}}{S_{11}(0)} \right)^{-2.5} \left(\frac{S_{33}}{S_{33}(0)} \right)^{-0.37} \left(\frac{S_{34}}{S_{34}(0)} \right)^{+0.8} \\ \times \left[1 + 3.47 \left(\frac{S_{17}}{S_{17}(0)} \right)^{+1.0} \left(\frac{\lambda_{e7}}{\lambda_{e7}(0)} \right)^{-1.0} \right] \\ \times \left(\frac{t_{\text{age}}}{4.7 \times 10^9 \text{yr}} \right)^{+1.4} \left(\frac{Z}{0.015} \right)^{+1.1}$$

Identified Important Reaction Flow Paths

Woodsley & Hoffman
(1992)



Sasaqui, Kajino, Mathews,
Otsuki & Nakamura
ApJ (2005), submitted.

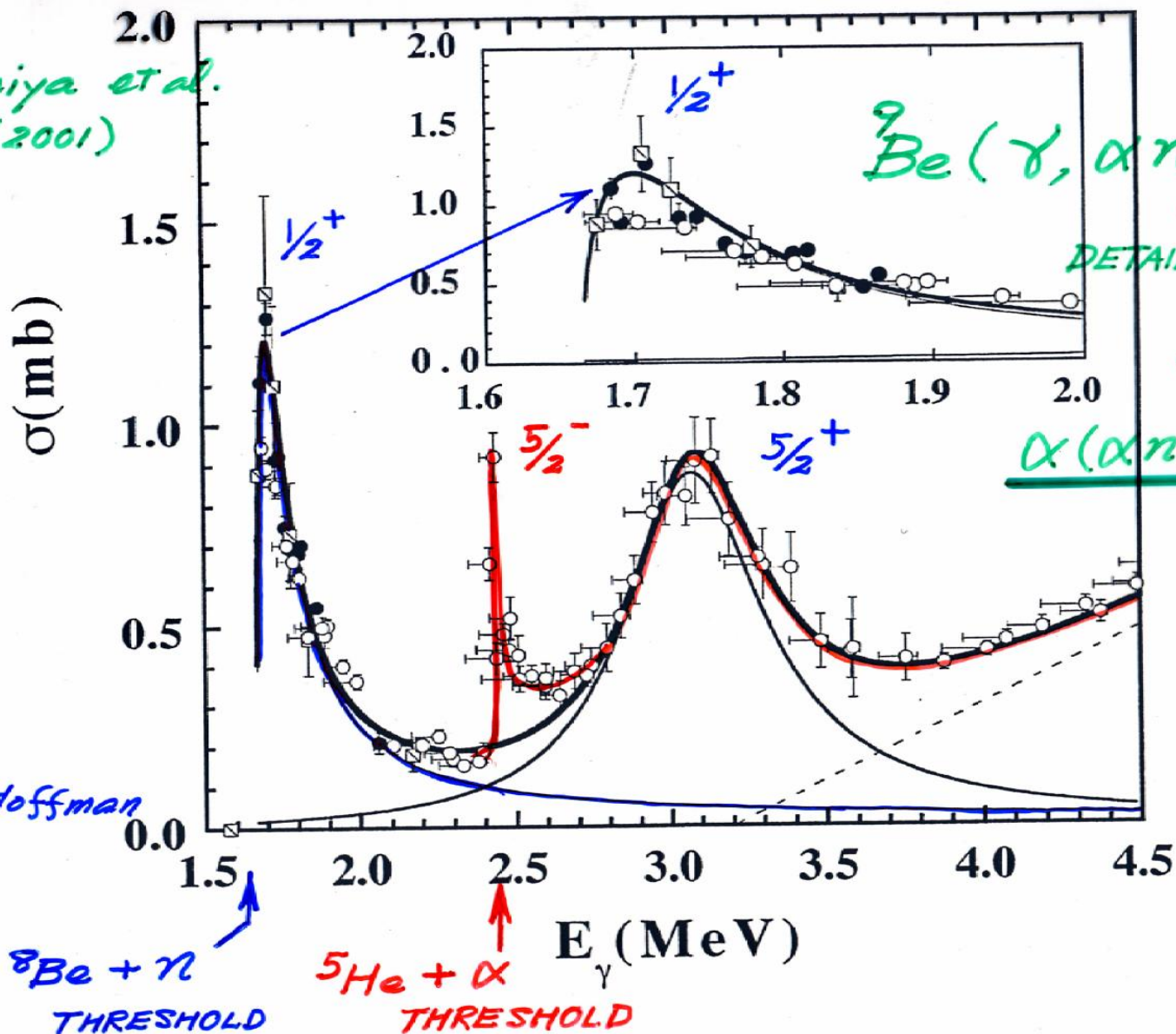
Terasawa, Sumiyoshi, Kajino,
Mathews & Tanihata
ApJ 562 (2001) 470.

(1) $\alpha(\alpha n, \gamma)^9\text{Be}(\alpha, n)^{12}\text{C}$

35%(1 σ)

EXP.

Utsunomiya et al.
PR C63 (2001)
018801



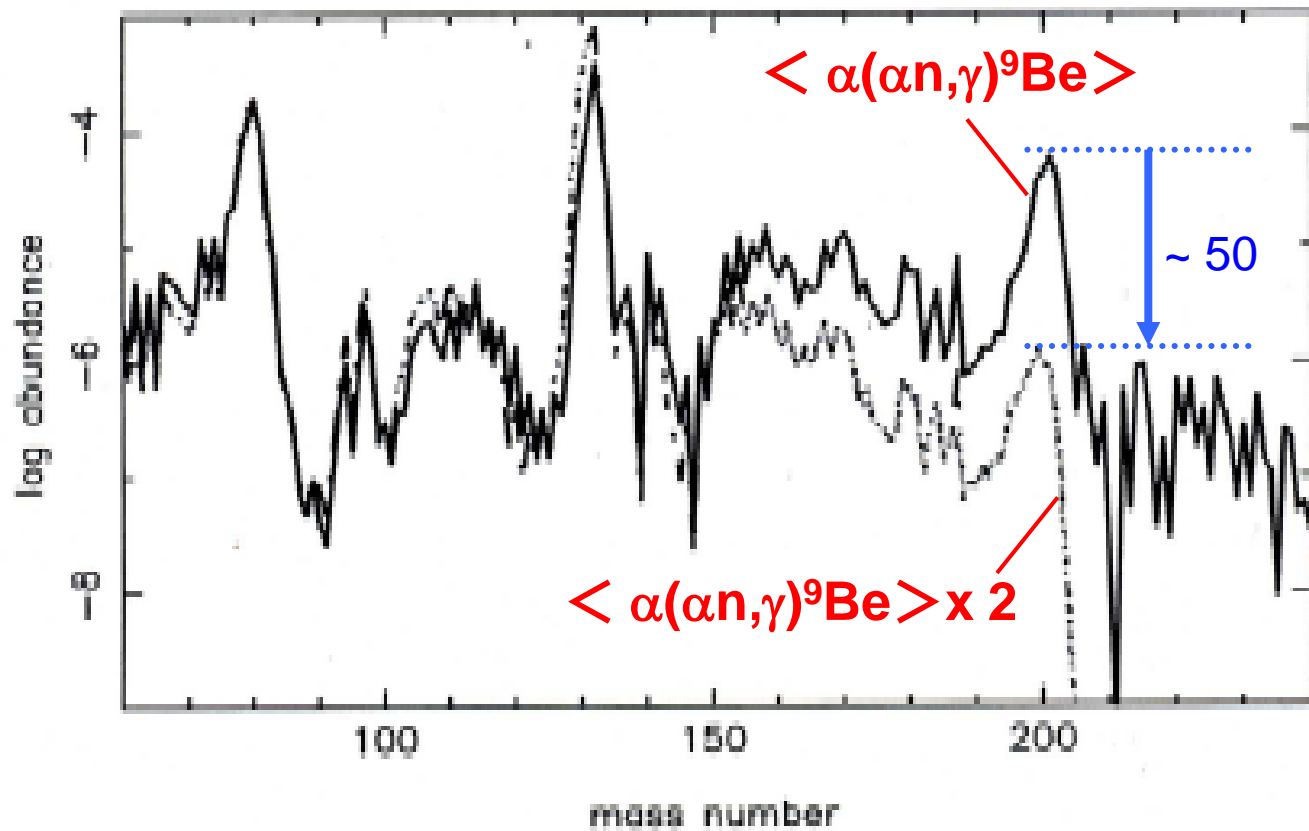
Woosley & Hoffman
(1992)

R-Process Sensitivity to Individual Reaction

Factor of 2 change of $\langle \alpha(\alpha n, \gamma)^9\text{Be} \rangle$ reaction rate

→ About factor 50 change in r-element yields !

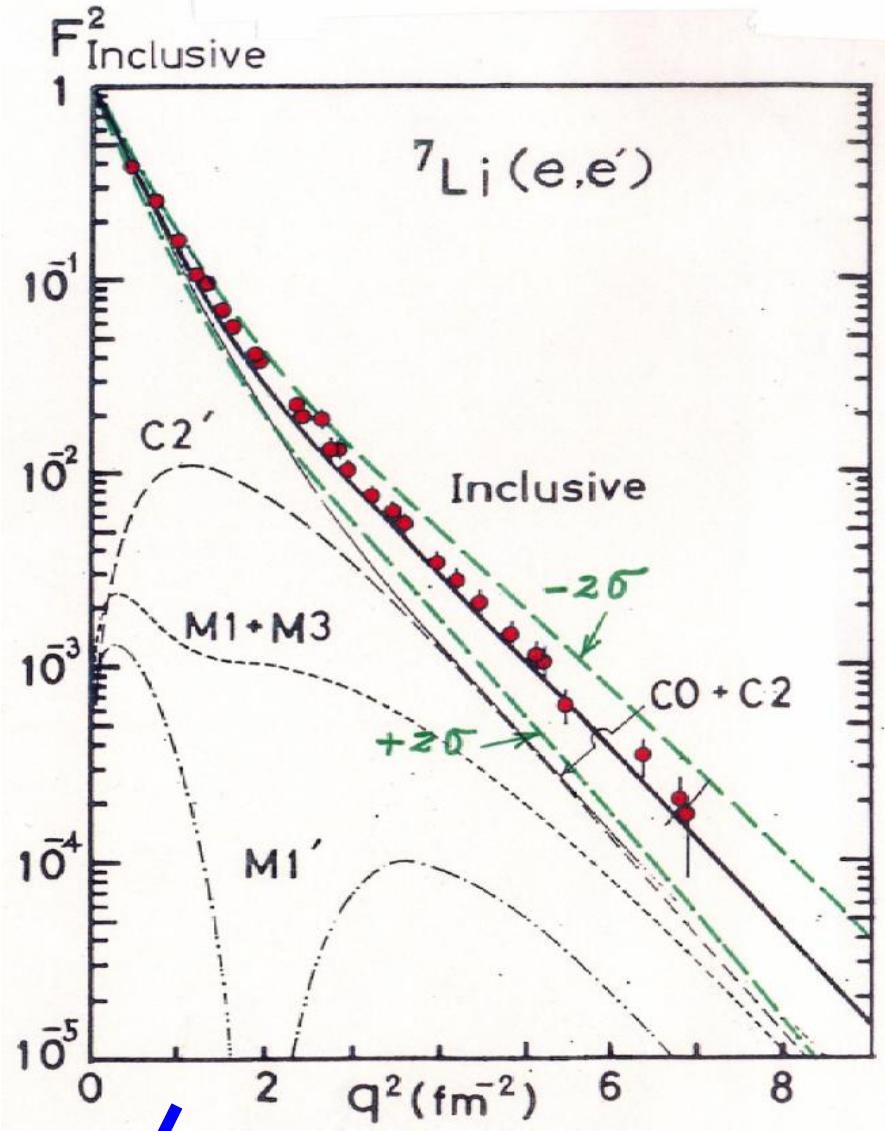
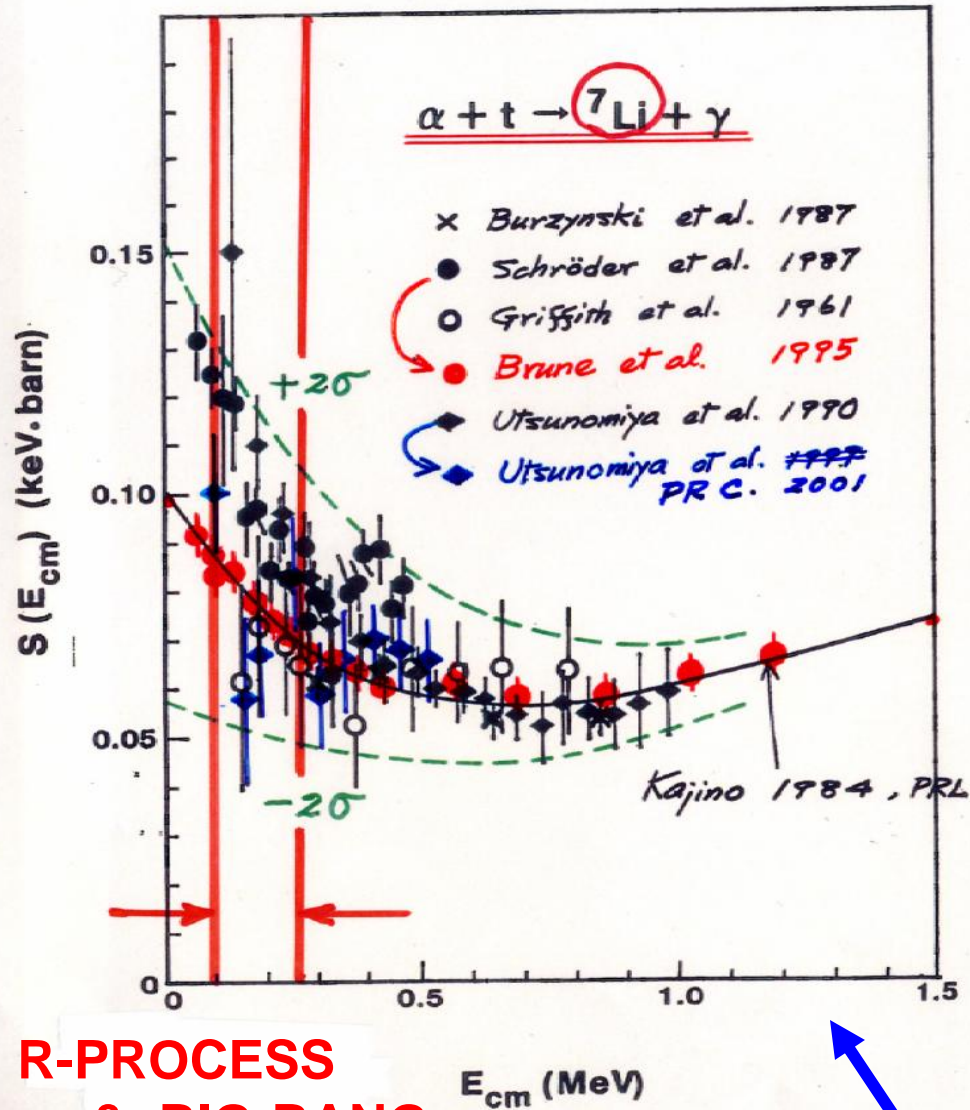
(slowly expanding v-wind model)



(2) $\alpha(t,\gamma)^7\text{Li}$

30%(1 σ)

Kajino et al. (2005)



R-PROCESS & BIG-BANG

REMOVING UNCERTAINTY !

(3) ${}^7\text{Li}(n,\gamma){}^8\text{Li}(\alpha,n){}^{11}\text{B}$

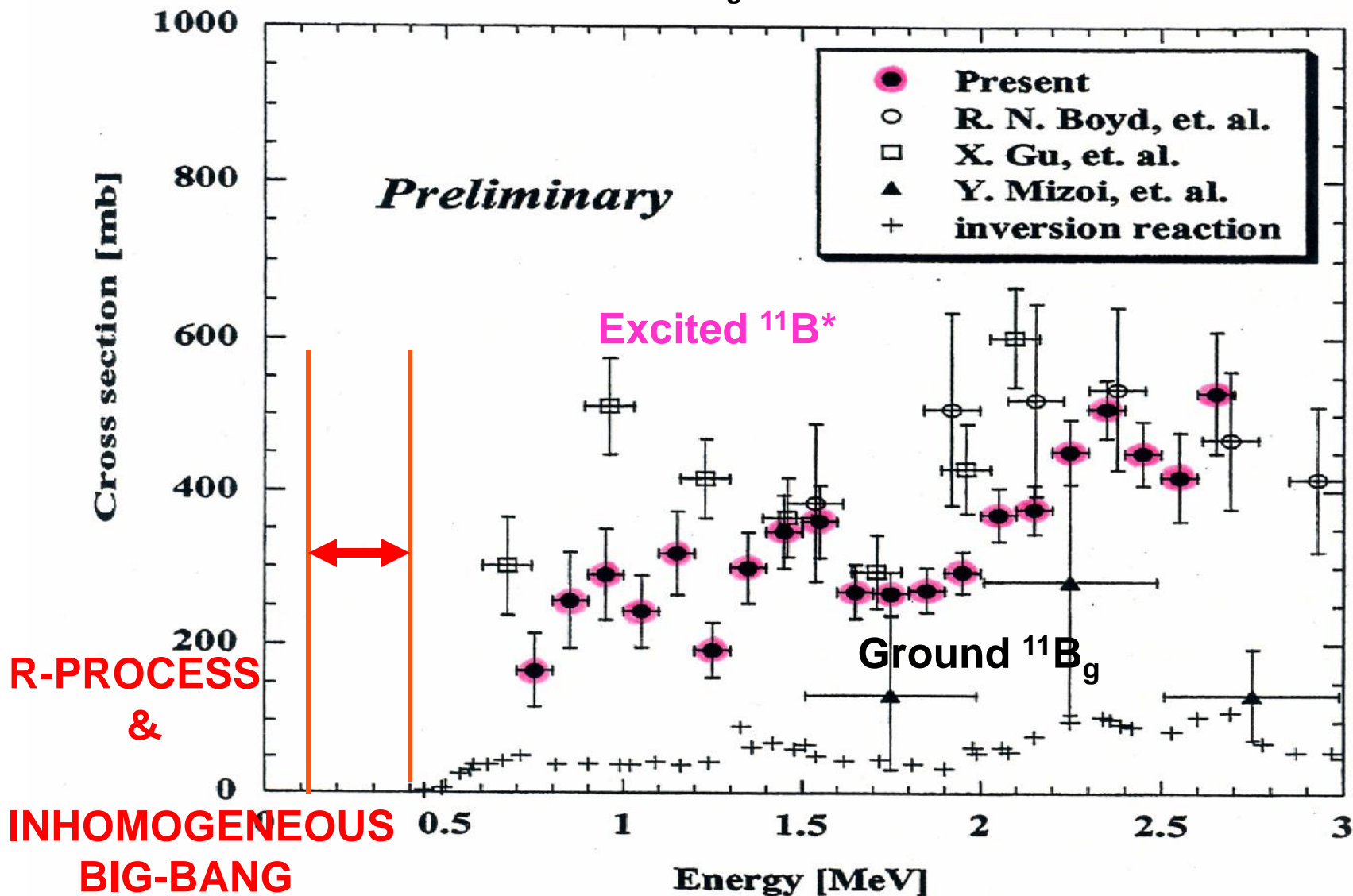
Factor 2 (1σ)

THEORY, unfinished.

H. Ishiyama et al. AIP Conf. Proc. 704 (2004) 453.

Yamamoto, Kubo, Ogawa & Kajino

${}^{11}\text{B}_g$ + several ${}^{11}\text{B}^*$ (EXCLUSIVE EXP.)



SENSITIVITY of Relevant Reactions to R-Process

Sasaqui, Kajino, Mathews, Otsuki & Nakamura, ApJ (2005) submitted.
 Otsuki, Tagoshi, Kajino & Wanajo, ApJ 533 (2000), 424.

$$Y_{0,r} + \delta Y_r = Y_{0,r} \{1 + 2\sigma\}^\alpha$$

(1) $\alpha(\alpha n, \gamma)^9\text{Be}$ $1\sigma = 35\%$ \longrightarrow $(Y_0 + \delta Y)/Y_0 = 0.35 \sim 11.2$

(2) $\alpha(t, \gamma)^7\text{Li}$ $1\sigma = 30\%$ \longrightarrow $0.27 \sim 13.2$

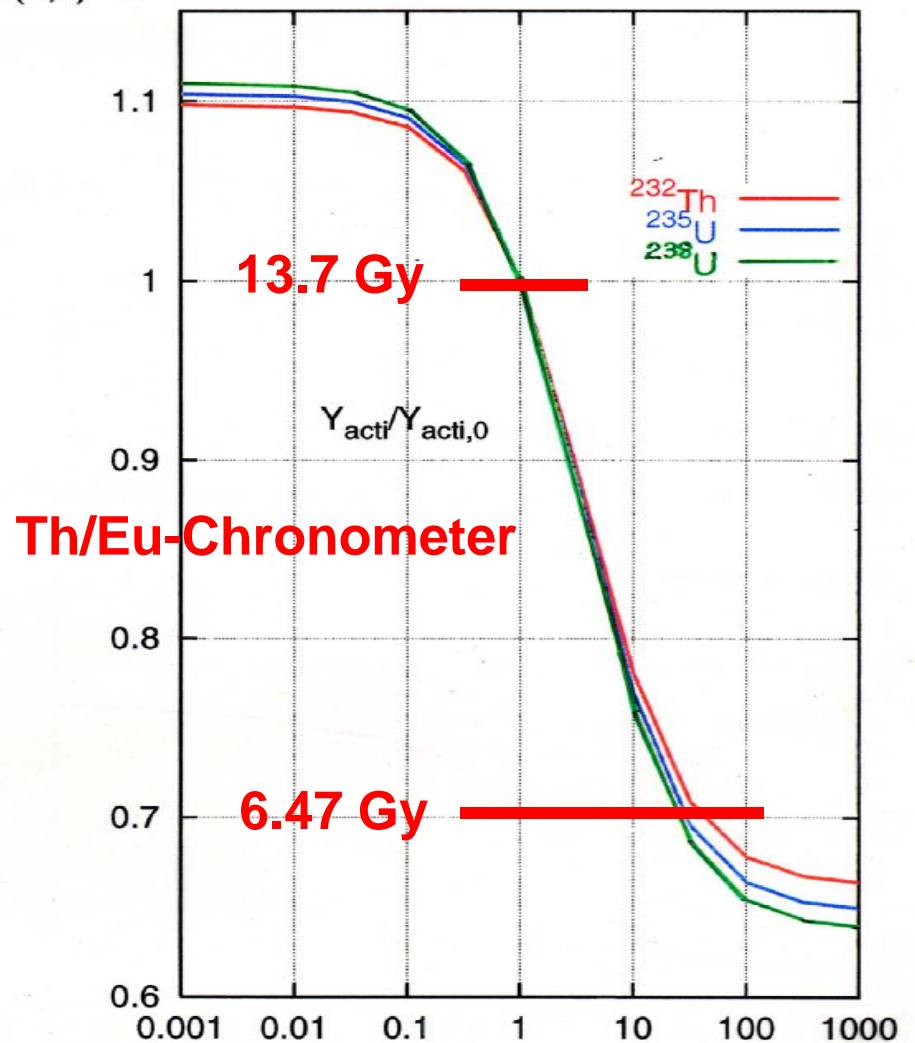
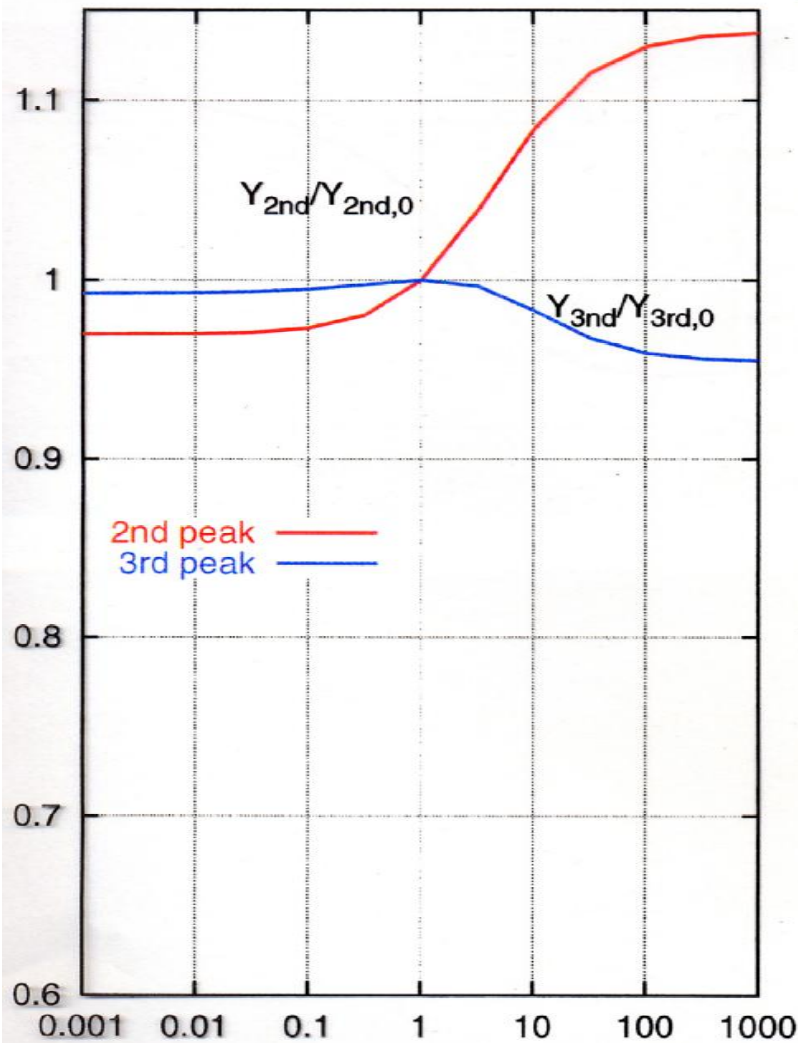
(3)(4) $^7\text{Li}(n, \gamma)^8\text{Li}(\alpha, n)^{11}\text{B}$ $1\sigma = 35\%, \times 2$ \longrightarrow $0.79 \sim 1.7$

(Th/U)=0.56-0.79

No.	reaction	sensitivity(α_i)					current	
		2nd peak	3rd peak	^{232}Th	^{235}U	^{238}U	importance	
(1)	$\alpha(\alpha n, \gamma)^9\text{Be}$	0.1823	-0.6546	-1.9423	-1.9819	-2.1006	0.3445	11.2222
(2)	$\alpha(t, \gamma)^7\text{Li}$	0.2874	-0.7474	-2.7125	-2.7857	-2.9583	0.2658	13.2353
(3)	$^7\text{Li}(n, \gamma)^8\text{Li}$	0.0465	-0.0917	-0.4296	-0.4436	-0.4729	0.7881	1.7163
(4)	$^8\text{Li}(\alpha, n)^{11}\text{B}$	0.0017	-0.0032	-0.0164	-0.0170	-0.0181	0.9882	1.0120

SENSITIVITY of ^{232}Th & $^{235,238}\text{U}$ to $^{18}\text{C}(\alpha,n)^{21}\text{O}$

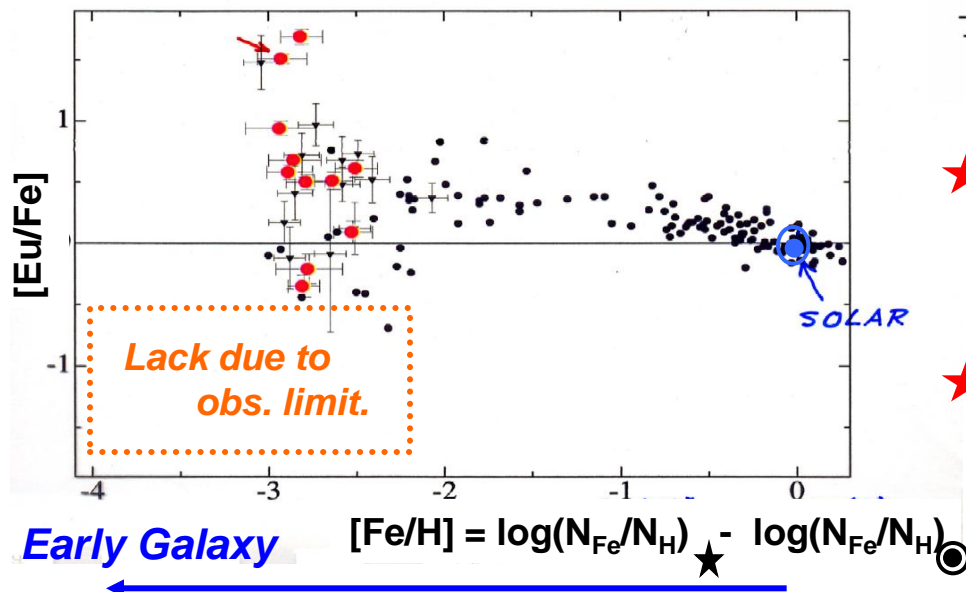
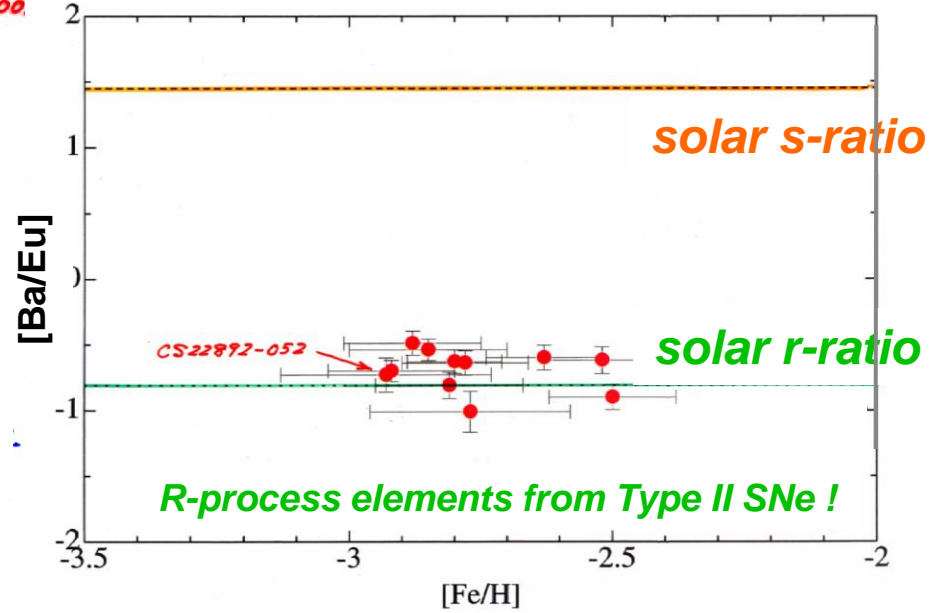
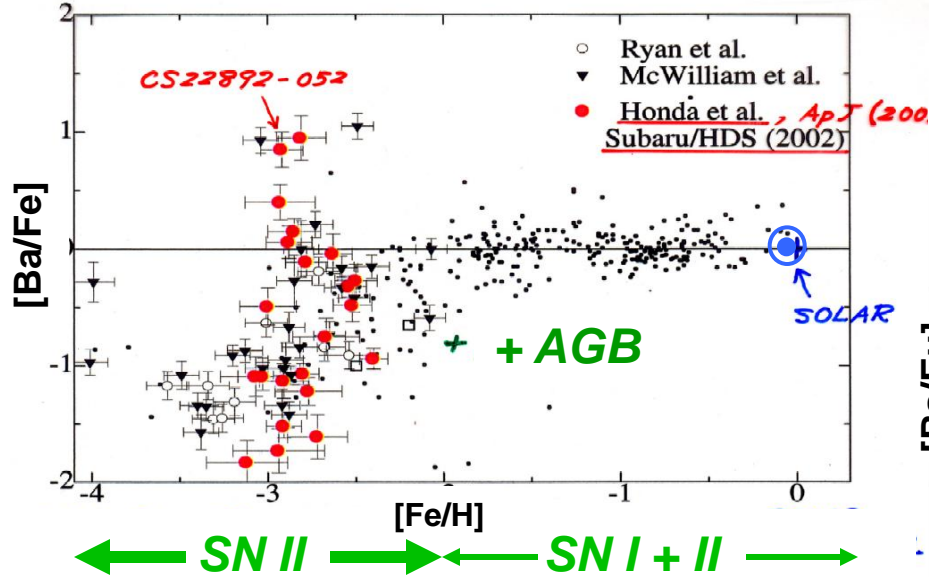
Sasaqui, Kajino, Mathews, Otsuki & Nakamura, ApJ (2005) submitted.



cross section ratio of $^{18}\text{C}(\alpha,n)^{21}\text{O}$

SUBARU Telescope HDS

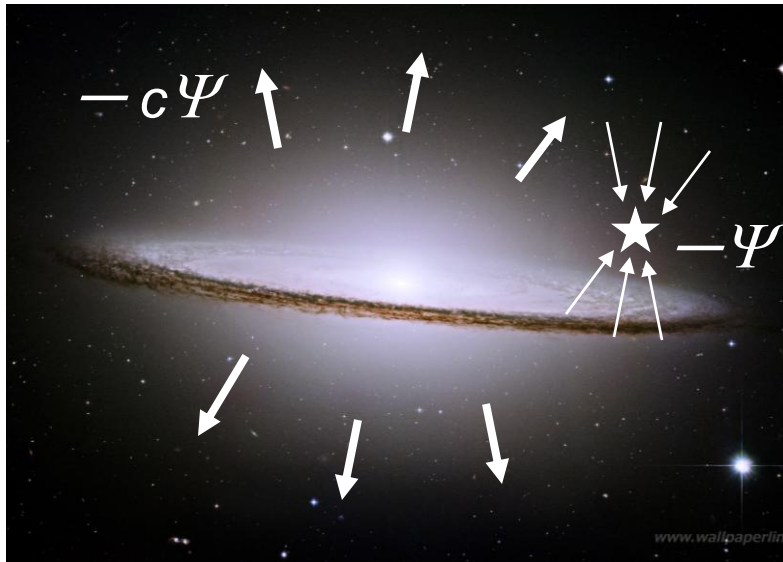
Honda, Aoki, Kajino et al.
 (SUBARU/HDS Collaboration),
 2004, ApJS 152, 113; 2004, ApJ 607, 474



- ★ Large abundance scatter at $[Fe/H] < -2$ is an evidence for INDIVIDUAL supernova episode.
- ★ Only Core-Collapse TYPE II SUPERNOVAE are the likely astrophysical sites of the R-Process !

Simple Galactic Chemical Evolution (GCE) Model

<http://www.kabegamilink.com/act/0704/03242.html>



Halo-Gas (M_G) and Stars ($M_{tot} - M_G$)

- Z_i = Mass Fraction of Nucleus- i
- y_i = Stellar Production Yield
- Ψ = Star Formation Rate
- ϕ = Galactic Cosmic Ray
- $c\Psi$ = Galactic Wind
- R = Returned Fraction $R = \sum R_i Z_i$

$$\left\{ \begin{array}{l} \frac{dM_{tot}}{dt} = -c\Psi \quad \text{--- (1)} \end{array} \right.$$

$$\frac{dM_G}{dt} = -(1-R+c)\Psi \quad \text{--- (2)}$$

$$\frac{d(M_G Z_i)}{dt} = y_i \Psi - (1-R_i+c)\Psi Z_i \quad \text{--- (3)}$$

$$\frac{d(M_G Z_L)}{dt} = \underbrace{y_L \Psi}_{\text{Stellar Production}} + \underbrace{\sum_j Z_j \left(\frac{A_L}{A_j}\right) \langle \sigma_{jL} \phi \rangle}_{\text{GCR production}} - (1-R_L+c)\Psi Z_L \quad \text{--- (4)}$$

Stellar Production

GCR production

結論

☆超新星Rプロセスはプライマリー過程。金属量に依らない。

- ・ 超金属欠乏星のすばる天文観測と合致。
- ・ 教科書き換え→
「Rプロセスは高温・高密度状態で数秒間に進行する爆発的元素合成過程であり、ばらばらの中性子と陽子から始まる。安定な鉄族元素を必要としない。」

☆重い原子核の性質 S_n 、 τ_β 、 $\sigma(n,\gamma)$ 、 $\sigma(v+A)$ ばかりでなく、軽い中性子過剰核の $\sigma(n,\gamma)$ 、 $\sigma(\alpha,n)$ が重要。

- ・ Th/U、Th/Eu など核宇宙時計の精度を保つために、原子核反応率をさらに5%程度の精度で決定する必要がある。

☆ビッグバン元素合成(プライマリー過程)と共通点多く、宇宙論モデルを強く制限する。